

Fun with Gamma, Beta, & error functions

Gamma

The Γ function is defined by integrals of the type

$$\Gamma(n) = \int_0^{\infty} x^{n-1} e^{-x} dx \quad (n > 0)$$

which, as you can imagine, appear A LOT!

This can't be evaluated analytically, hence we just call it $\Gamma(n)$

let's evaluate $\Gamma(n+1)$

$$\Gamma(n+1) = \int_0^{\infty} x^{(n+1)-1} e^{-x} dx = \int_0^{\infty} x^n e^{-x} dx$$

you should be very tempted to integrate by parts

$$\Gamma(n+1) = \left[-x^n e^{-x} \right]_0^{\infty} + \int_0^{\infty} n x^{n-1} e^{-x} dx$$

verify that $\boxed{\Gamma(n+1) = n \Gamma(n)}$

hint e^x always beats x^n as $x \rightarrow \infty$

notice that $\Gamma(1) = 1$

$$\Rightarrow \Gamma(2) = 1 = \Gamma(1+1)$$

$$\Rightarrow \Gamma(2+1) = 2\Gamma(2)$$

$$\Rightarrow \Gamma(3+1) = 3\Gamma(3) = 3 \times 2\Gamma(2) \\ = 3 \times 2 \times 1 = 3!$$

by induction

$$\boxed{\Gamma(n+1) = n!}$$

$\Gamma(n+1)$ gives us a new definition of factorial even for - values of n

let $x = y^2$ then $\Gamma(n) = 2 \int_0^{\infty} y^{2n-1} e^{-y^2} dy$
 $dx = 2y dy$

let $n = \frac{1}{2} \Rightarrow \Gamma(\frac{1}{2}) = 2 \int_0^{\infty} y^0 e^{-y^2} dy$

but that's a gaussian

$$\int_{-\infty}^{\infty} e^{-y^2} dy = \sqrt{\pi}$$

$$\Rightarrow \Gamma(\frac{1}{2}) = \sqrt{\pi}$$

from $\Gamma(n+1) = n\Gamma(n)$ you get all sorts of values!

Here's one last crazy useful idea,

put the x^n into an exponential by $e^{n \ln x} = x^n$

then
$$n! = \int_0^{\infty} e^{n \ln x - x} dx$$

now change variables $x = n + y$ or $y = x - n$

this isn't an obvious thing to do but

$dx = dy$;

$$\ln x = \ln n + \ln\left(1 + \frac{y}{n}\right)$$

for large n we can Taylor expand $\ln\left(1 + \frac{y}{n}\right)$

$$\ln\left(1 + \frac{y}{n}\right) \approx \frac{y}{n} - \frac{y^2}{2n^2} + \dots$$

plugging back into $\Gamma(n)$

$$n! = \int_{-n}^{\infty} e^{n \left[\ln n + \frac{y}{n} - \frac{y^2}{2n^2} + \dots \right] - n - y} dy$$

$$\approx e^{n \ln n - n} \int_{-\infty}^{\infty} e^{-\frac{y^2}{2n}} dy = e^{n \ln n - n} \sqrt{2\pi n}$$

because we let
n become very large

$$\therefore \boxed{n! \approx n^n e^{-n} \sqrt{2\pi n}}$$

This is called
Stirling's approximation &
is VERY useful!

Beta

$$\begin{aligned} B(m, n) &= \int_0^1 x^{m-1} (1-x)^{n-1} dx \\ &= \int_0^\infty \frac{y^{n-1} dy}{(1+y)^{m+n}} \end{aligned}$$

$B(m, n)$ isn't terribly useful but it is related to $\Gamma(n)$ by

$$B(m, n) = \frac{\Gamma(m) \Gamma(n)}{\Gamma(m+n)}$$

this form actually shows up in String Theory!

Error function

$$\begin{aligned} \operatorname{erf}(x) &= \frac{2}{\sqrt{\pi}} \int_0^x e^{-u^2} du \\ &= 1 - \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-u^2} du \end{aligned}$$

This guy shows up in probability. one generally looks up ~~is~~ its values or numerically compute them

$\Gamma(n)$ $\hat{=}$ spheres

we saw $\Gamma(\frac{1}{2}) = \sqrt{\pi}$ $\hat{=}$ prompts us to wonder "where the heck are circles?!"

let's find the volume of a sphere.

consider 3D (x, y, z)

a sphere is the surface defined by

$$x^2 + y^2 + z^2 = R^2$$

when we say "sphere" we always mean the 2D surface.

the "Ball" is the surface along with the interior.

The volume of a Ball is given by

$$V_S = \int_S dV = \int r^2 \sin \theta dr d\theta d\phi = \frac{4}{3} \pi R^3$$

what about arbitrary spheres of any dimension?

again we want to solve

$$V_{HS} = \int dV_n \quad \text{where}$$

$$x_1^2 + x_2^2 + \dots + x_n^2 = R^2$$

is an n -dimensional Sphere (the surface is $(n-1)$ dimensional)

if we knew dV_n we could just find V_{HS}

1D
 $dV_1 = dr$

2D
 $dV_2 = r dr d\theta$

3D
 $dV_3 = r^2 \sin\theta dr d\theta d\phi$

the angular parts of dV is often denoted $d\Omega$

where Ω is a solid angle

e.g. $dV_3 = r^2 dr d\Omega_2$

Now, notice what $\int d\Omega_2$ gives us.

$$\int d\Omega = \text{surface area of our sphere}$$

in arbitrary dimensions we expect

$$\boxed{dV_n = r^{n-1} dr d\Omega_{n-1}}$$

Now, if we could somehow figure out the surface area of a sphere then we could easily find the volume by integrating

$$\int r^{n-1} dr \quad \text{is multiplying by } S_{n-1} = \int d\Omega_{n-1}$$

Here's the insanely cool trick!

compute $\int_{-\infty}^{\infty} e^{-x^2} dx = I$

we previously found this by squaring
then going to polar coords & solving

$$I^2 = \iint e^{-x^2-y^2} dx dy$$

$$\Rightarrow \int_0^{2\pi} \int_0^{\infty} \underbrace{r dr \frac{d\theta}{d\Omega}}_{d\Omega} e^{-r^2} = \pi$$

$$\Rightarrow I = \sqrt{\pi}$$

were more than welcome to compute I^3

$$\begin{aligned} I^3 &= \iiint e^{-(x^2+y^2+z^2)} dx dy dz \\ &= (\sqrt{\pi})^3 \end{aligned}$$

or even $I^n = (\sqrt{\pi})^n = \int e^{-(x_1^2+x_2^2+\dots+x_n^2)} dx_1 dx_2 dx_3 \dots dx_n$

But in polar coordinates we get that
same angular part $d\Omega$

$$I^n = \pi^{n/2} = \int r^{n-1} e^{-r^2} dr d\Omega$$

Oh my god look at that radial integral!

does it look similar to something you've seen?!

$$\Gamma(n) = 2 \int_0^{\infty} y^{2n-1} e^{-y^2} dy$$

so that our radial integral is "simply"

$$\frac{1}{2} \Gamma\left(\frac{n}{2}\right)$$

all ~~to get~~ together

$$\begin{aligned} I^n &= \pi^{n/2} = \frac{1}{2} \Gamma\left(\frac{n}{2}\right) \int d\Omega \\ &= \frac{1}{2} \Gamma\left(\frac{n}{2}\right) (\text{Surface area}) \end{aligned}$$

solving for the surface area

$$S_n = \frac{2 \pi^{n/2}}{\Gamma\left(\frac{n}{2}\right)}$$

then the

$$\text{Volume} = \int_0^R S_n r^{n-1} dr = \frac{2 R^n \pi^{n/2}}{n \Gamma\left(\frac{n}{2}\right)}$$

is the volume of an n -dimensional sphere. Wow!